

# The Battery's New Clothes: Structure and Li-ion dynamics of the Crystalline Lithium-Rich Anti-Perovskites $\text{Li}_3\text{OCl}$ and $\text{Li}_2(\text{OH})\text{Cl}$

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Lithium-rich anti-perovskites (LiRAPs) might belong to the next-generation of promising solid electrolytes for all-solid-state batteries. Their 3D  $\text{Li}^+$  migration channels are supposed to be beneficial for rapid Li ion dynamics [1, 2]. Their reliable preparation is, however, fraught with difficulties, especially when glassy samples have to be prepared [3]. Thermodynamic instabilities, sensitivity to moisture, and impurity phases often hinder the systematic study of ion transport properties in this class of materials.

Here, we prepared phase pure LiRAPs by hydrothermal synthesis with and without subsequent annealing steps. We investigated both structure and ionic conductivity and diffusivity of  $\text{Li}_3\text{OCl}$  and compared the results to those of  $\text{Li}_2(\text{OH})\text{Cl}$ .  $\text{Li}_3\text{OCl}$  transforms into  $\text{Li}_2(\text{OH})\text{Cl}$  if proton sources are present. We applied a variety of methods, such as X-ray diffraction (XRD), neutron diffraction, conductivity spectroscopy,  $^7\text{Li}$  NMR line shape measurements and NMR relaxometry experiments, to study the crystal chemistry as well as short-range ion dynamics and long-range ion transport. As revealed by XRD,  $\text{Li}_2(\text{OH})\text{Cl}$  undergoes a reversible phase transition at 313 K. This transition is associated with a jump in ionic conductivity by three orders of magnitude yielding  $\sigma(313 \text{ K}) = 5.1 \times 10^{-6} \text{ S cm}^{-1}$ . Such behaviour was, however, not seen for phase pure H-free  $\text{Li}_3\text{OCl}$ , which is characterized by the following parameters ( $\sigma_{\text{bulk}}(293 \text{ K}) = 0.84 \times 10^{-6} \text{ S cm}^{-1}$ ,  $0.63(1) \text{ eV}$ ;  $\sigma_{\text{grain boundary}}(293 \text{ K}) = 0.28 \times 10^{-6} \text{ S cm}^{-1}$ ,  $E_a = 0.62(2) \text{ eV}$ ). The steep increase in ion conductivity of  $\text{Li}_2(\text{OH})\text{Cl}$  is also corroborated by  $^7\text{Li}$  NMR line widths ( $\Delta_{\text{NMR}}$ ) measurements. Significantly narrowed NMR lines show up above 313 K ( $\Delta_{\text{NMR}} = 1100 \text{ Hz}$ ) due to rapid  $\text{Li}^+$  jump processes in LiRAPs that successfully average dipolar interactions giving rise to broadened lines at much lower temperature ( $\Delta_{\text{NMR}}(273 \text{ K}) = 8700 \text{ Hz}$ ). Our measurements suggest that besides  $\text{Li}^+$  also other charge carriers contribute to the overall conductivity of  $\text{Li}_2(\text{OH})\text{Cl}$ . While long-range  $\text{Li}^+$  transport in  $\text{Li}_2(\text{OH})\text{Cl}$  is governed by  $0.57(1) \text{ eV}$ ,  $^7\text{Li}$  NMR relaxometry senses local ion jumps with an activation energy of  $0.36(1) \text{ eV}$ . The latter value includes the effect of correlated motion on NMR activation energies measured in the low-temperature regime of the corresponding NMR spin-lattice relaxation rate peak, see Ref. [4]. Variable-temperature measurements of the real part of the complex resistivity ( $\rho' = M''/\omega$ ), see Ref. [5] for another example, revealed an asymmetric peak with activation energies of  $0.38(2) \text{ eV}$  and  $0.55(2) \text{ eV}$ . These results clearly show complex dynamics in  $\text{Li}_3\text{OCl}$  presumably including both translational and (interrelated) rotational motions.

## References:

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